

# An open-source perspective at CFD-DEM for particle-laden flows

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## Polytechnique Montréal



10000STUDENTS47,500GRADUATES<br/>SINCE 1873,2000DEGREES<br/>AWARDED<br/>LAST YEAR

300 PROFESSORS

\$250M ANNUAL BUDGET



## CHAOS Lab

Chem. Eng. High-performance Automatization, Optimization and Simulation

Group					
Post-doc	1				
PhD	8				
Master	5				







## Particle-laden flow

Industrial applications of interest to us



## Solid-fluid contactor: Fluidized-bed

## Fluidized bed contactor

- Uniform fluid jet
- Strong particle-fluid mixing
- Many regimes
- Geldart B particles



## Solid-fluid contactor: Spouted-bed

## Spouted-bed contactor

- Air jet in a small slit
- Contact in upward movement of fluid
- Geldart D particles





## Pneumatic conveying





# Novel processes: LPBF 3D printing





# Summary

Quick recap on high-order methods

- Lethe: open-source high-order CFD/DEM/CFD-DEM
- Illustration using two turbulent benchmarks
- High-order Unresolved CFD-DEM
- VANS equation formulation
- Void fraction projection scheme
- Interesting validation cases

Conclusions and future work





# High-order methods for CFD

A quick introduction



## Order of a scheme?

The order characterizes how the error will evolve as the mesh (or time step) is refined (or coarsened)

l.e. :





## Traditional CFD

The traditional state of the art in commercial and open source software is mostly centered around 2<sup>nd</sup> order accurate scheme

Commercial: • Fluent • Star-CCM+ Open-source • OpenFOAM • SU2

Why stop there?





# High-order and FEM

In the case of FEM using high-order methods entails using high order polynomial interpolations. This can be done using either: • Continuous Galerkin - CG

• Discontinuous Galerkin - DG





# Efficiency of high-order methods

Same accuracy in shorter solution time

Depends on:

Geometry

• Reynolds number





## Lethe





# Methors

# Solid open-source foundations

- Based on the deal.II framework
- Tested (>340 tests run daily)
- Documented (>50 examples)
- Used in 5 countries
- (Canada, Germany, UK, Australia, USA)

What is the FEM formulation used in Lethe for the Navier-Stokes equations?

### **Incompressible Flow**



## Continuity

$$\int_{\Omega} \nabla \cdot \boldsymbol{u} q \mathrm{d}\Omega + \underbrace{\sum_{K} \int_{\Omega_{k}} \mathcal{SR} \cdot (\tau_{\boldsymbol{u}} \nabla q) \mathrm{d}\Omega_{k}}_{\mathrm{PSPG}} = 0$$

### Momentum

$$\int_{\Omega} \left( \frac{\partial \boldsymbol{u}}{\partial t} + \boldsymbol{u} \cdot \nabla \boldsymbol{u} \right) \cdot \boldsymbol{v} d\Omega + \int_{\Omega} \boldsymbol{\nu} \left( \nabla \boldsymbol{u} : \nabla \boldsymbol{v} \right) d\Omega$$
$$- \int_{\Omega} \left( p \nabla \cdot \boldsymbol{v} \right) d\Omega + \sum_{K} \int_{\Omega_{k}} S\mathcal{R} \cdot \left( \tau_{u} \boldsymbol{u} \cdot \nabla \boldsymbol{v} \right) d\Omega_{k} = 0$$
SUPG

## Comments

## Stabilized formulation:

- Allows QnQn elements, circumvents LBB conditions
- Implicit LES

Non-linear problem solution strategy
 Newton-Raphson method with dynamic recalculation of the Jacobian matrix

### Linear solver strategy

Monolithic solver using AMG, ILU or GMG preconditioning

## High order

• High order methods move the numerical dissipation to higher wave number. Closer to where the real viscosity is acting







# A simple turbulent benchmark

Taylor-Green vortex at Re=1600

Wang, Zhijian J., et al. "High-order CFD methods: current status and perspective." International Journal for Numerical Methods in Fluids (2013) – Test 3.5

Blais et al. "Lethe: An open-source parallel high-order adaptative CFD solver for incompressible flows." *Software X* 12 (2020)



## Taylor-Green test case

Unsteady flow of decaying vortices

In 2D:

• Exponential decay of the vortices

In 3D:

- Generation of a fully 3D flow patterns
- Turbulent cascade that generate smaller flow structures which then dissipate

The TGV test case is a great benchmark for DNS and ILES

$$u = V_0 \sin\left(\frac{x}{L}\right) \cos\left(\frac{y}{L}\right) \cos\left(\frac{z}{L}\right)$$
$$v = -V_0 \cos\left(\frac{x}{L}\right) \sin\left(\frac{y}{L}\right) \cos\left(\frac{z}{L}\right)$$
$$w = 0$$
$$p = p_0 + \frac{\rho V_0^2}{16} \left(\cos\left(\frac{2x}{L} + \cos\left(\frac{2y}{L}\right)\right)\right) \left(\cos\left(\frac{2z}{L}\right) + 2\right)$$



# Velocity

Cascade from large turbulent structures to smaller ones

Small structures dissipate the energy





# Energy dissipation - Enstrophy

$$\left[\frac{\partial}{\partial t}\left(\frac{u^2}{2}\right) = -\nabla \cdot \left[\left(\frac{u^2}{2} + \frac{p}{\rho}\right)\mathbf{u} + \nu\left(\nabla \times \mathbf{u}\right) \times \mathbf{u}\right] - \nu\left(\nabla \times \mathbf{u}\right)^2$$

Integrating on a closed or periodic domain, we obtain:

$$\frac{\partial}{\partial t} \int \left(\frac{u^2}{2}\right) dV = \int -\nu \left(\nabla \times \mathbf{u}\right)^2 dV$$

In absence of numerical dissipation, the integral of the enstrophy should be equal to the decay of kinetic energy



 $Q1-Q1 - 256x^3$ 



### Bether

 $Q2-Q2-128^{3}$ 

Dramatic change

Good agreement for enstrophy

<u>67.9M DOFs</u>





Q3-Q3 -128<sup>3</sup>







Q3-Q3 -64<sup>3</sup>



#### Methom



## Enstrophy over time





# High-order Unresolved CFD-DEM

Model formulation



## Volume-Averaged Navier-Stokes equations

$$\frac{\partial (\epsilon_f)}{\partial t} + \nabla \cdot (\epsilon_f \boldsymbol{u}_f) = 0$$
$$\frac{\partial (\epsilon_f \boldsymbol{u}_f)}{\partial t} + \nabla \cdot (\epsilon_f \boldsymbol{u}_f \otimes \boldsymbol{u}_f) = -\frac{\epsilon_f}{\rho_f} \nabla p + \epsilon_f \nabla \cdot (\boldsymbol{\tau}_f) + \frac{\boldsymbol{F_{pf}^A}}{\rho_f}$$

We discretize the VANS equations using stabilized FEM



### Continuity

$$\int_{\Omega} \left( \frac{\partial \left( \epsilon_{f} \right)}{\partial t} + \epsilon_{f} \nabla \cdot \boldsymbol{u} + \boldsymbol{u} \nabla \epsilon_{f} \right) q d\Omega + \underbrace{S\mathcal{R} \cdot \left( \tau_{\boldsymbol{u}} \nabla q \right) d\Omega_{k}}_{\text{PSPG}} = 0$$

### Momentum

$$\int_{\Omega} \rho_f \left( \epsilon_f \frac{\partial \boldsymbol{u}}{\partial t} + \epsilon_f \boldsymbol{u} \cdot \nabla \boldsymbol{u} \right) \cdot \boldsymbol{v} d\Omega + \int_{\Omega} \left( \epsilon_f \boldsymbol{\nu} \left( \nabla \boldsymbol{u} : \nabla \boldsymbol{v} \right) \right) + \left( \boldsymbol{\nu} \nabla \boldsymbol{u} \nabla \epsilon_f \cdot \boldsymbol{v} \right) d\Omega$$
$$- \frac{1}{\rho_f} \int_{\Omega} \left( \epsilon_f \boldsymbol{p} \cdot \nabla \boldsymbol{v} + \boldsymbol{p} \nabla \epsilon_f \cdot \boldsymbol{v} \right) d\Omega + \frac{1}{\rho_f} \int_{\Omega} \frac{\boldsymbol{F_{pf}}}{V_{\Omega}} \cdot \boldsymbol{v} d\Omega d\Omega$$
$$+ \sum_K \int_{\Omega_k} \mathcal{SR} \cdot \left( \tau_u \boldsymbol{u} \cdot \nabla \boldsymbol{v} \right) d\Omega_k + \sum_K \int_{\Omega_k} \gamma \left( \frac{\partial \epsilon_f}{\partial t} + \nabla \cdot \left( \epsilon_f \boldsymbol{u} \right) \right) \left( \nabla \cdot \boldsymbol{v} \right) d\Omega_k = 0$$
$$\text{SUPG}$$

## Multiple stabilization term

$$S\mathcal{R} = \frac{\partial \left(\epsilon_{f} \boldsymbol{u}_{f}\right)}{\partial t} + \nabla \cdot \left(\epsilon_{f} \boldsymbol{u}_{f} \otimes \boldsymbol{u}_{f}\right) + \frac{\epsilon_{f}}{\rho_{f}} \nabla p - \epsilon_{f} \nu \nabla^{2} \boldsymbol{u} - \frac{\boldsymbol{F}_{\boldsymbol{p}\boldsymbol{f}}^{\boldsymbol{A}}}{\rho_{f}}$$

**PSPG: Enables equal-order elements**  $\sum_{K} \int_{\Omega_{k}} S\mathcal{R} \cdot (\tau_{u} \nabla q) d\Omega_{k}$ 

**SUPG: Implicit LES model**  $\sum_{K} \int_{\Omega_{k}} S\mathcal{R} \cdot (\tau_{u} \boldsymbol{u} \cdot \nabla \boldsymbol{v}) d\Omega_{k}$ 

**Grad-div: Enhances mass conservation**  $\sum_{K} \int_{\Omega_{k}} \gamma \left( \frac{\partial \epsilon_{f}}{\partial t} + \nabla \cdot (\epsilon_{f} \boldsymbol{u}) \right) (\nabla \cdot \boldsymbol{v}) d\Omega_{k}$ 

## Simple example

- Flow through a pipe
  Slip on the walls
- Constant velocity inletJump in the void fraction
- oump in the volu nuc

$$\epsilon_f = \epsilon_f(\boldsymbol{x})$$

Mass must be conserved, so velocity must increase if the porosity decreases



## Impact of stabilization

SUPG + PSPG leads to strong oscillation when the void fraction varies Additional Grad-div stabilization in the momentum equation remedies this issue



## Discrete Element Method

Newton second's law solved for each particle

$$m_i \frac{\partial \boldsymbol{u}_{p,i}}{\partial t} = \sum_{j \in C_i} \boldsymbol{f}_{c,ij} + \sum_{w} \boldsymbol{f}_{c,iw} + m_i \boldsymbol{g} + \boldsymbol{f}_{\boldsymbol{fp},i}$$

Contact forces are calculated by allowing particle overlap

Overlap is decomposed into two directions

• Normal

• Tangential





## Contact detection

Binned list • O(n<sub>p</sub>) scaling

- Superposed by a Verlet list
- Algorithmic decisions
- List is refreshed every n iterations dynamically
- CFD mesh is used as contact detection structure
- Domain decomposition is used to parallelize the calculations



# Parallelization and load-balancing

- Underlying grid can be decomposed
- As particle move, load dynamically evolve
- Parallel decomposition must thus follow





# Load-balancing



Methor

## Void fraction calculation

## Challenges:

- Continuity in time
- Ideally as independent of the mesh as possible

# Solution is a spherical quadrature centered void fraction calculation







## **Unresolved CFD-DEM model**

### Fluid

Volume-Averaged Navier-Stokes (VANS form A)

$$\frac{\partial (\epsilon_f)}{\partial t} + \nabla \cdot (\epsilon_f \boldsymbol{u}_f) = 0$$
$$\frac{(\epsilon_f \boldsymbol{u}_f)}{\partial t} + \nabla \cdot (\epsilon_f \boldsymbol{u}_f \otimes \boldsymbol{u}_f) = -\frac{\epsilon_f}{\rho_f} \nabla p + \frac{\epsilon_f}{\rho_f} \nabla \cdot (\boldsymbol{\tau}_f) + \frac{\boldsymbol{F}_{\boldsymbol{p}f}^{\boldsymbol{A}}}{\rho_f}$$

### Solid particles

Newton's second law

$$m_i \frac{\partial \boldsymbol{u}_{p,i}}{\partial t} = \sum_{j \in C_i} \boldsymbol{f}_{c,ij} + \sum_{w} \boldsymbol{f}_{c,iw} + m_i \boldsymbol{g} + \boldsymbol{f}_{\boldsymbol{fp},i}$$

### Hydrodynamic forces

 $\boldsymbol{F}_{pf} = \frac{1}{\Delta V} \sum_{i}^{n_{p}} \boldsymbol{f}_{pf,i}$  $\boldsymbol{f}_{pf,i} = \boldsymbol{f}_{d,i} + \boldsymbol{f}_{\nabla p,i} + \boldsymbol{f}_{\nabla \cdot \boldsymbol{\tau},i} + \boldsymbol{f}_{\text{Saff},i}$ 

Pressure gradient

Viscous stress

- Drag
- Lift forces



## Case 1: Spouted bed

### Unresolved CFD-DEM

El Geitani, T., Golshan, S., & Blais, B. (2023). Toward High-Order CFD-DEM: Development and Validation. *Industrial & Engineering Chemistry Research*, *62*(2), 1141-1159.



## Spouted bed contactor

## Physical parameters

$$\circ$$
 Glass beads ( $ho_s=2500rac{kg}{m^3}$  and  $d_p=2.5mm$ 

- Approx 180k particles
- Air
- 30s of flow simulated

El Geitani, T., Golshan, S., & Blais, B. (2023). Toward High-Order CFD-DEM: Development and Validation. *Industrial & Engineering Chemistry Research*, *62*(2), 1141-1159.





## Validation strategy





## Q1 results



1.2

Experiment

Simulation

.

## Q2 results



Experiment

•

## A more direct comparison







# Case 2: Solid-liquid mixing drum

A simple benchmark for unresolved CFD-DEM

El Geitani, Toni, and Bruno Blais. "Solid-liquid rotary kilns: An experimental and CFD-DEM study." Powder Technology 430 (2023): 119008.



# Solid-liquid mixing drum





## Experimental setup



Besides, Froude number (Fr), the fill level as well as the particles' friction coefficient affect the regime change.

$$Fr = \frac{\omega^2 R}{g}$$



## Materials used

### Particles' Types



(a) 3 mm glass beads, (b) 5 mm glass beads and (c) 5.95 mm Acrylonitrile butadiene styrene (ABS) particles

### Fluids' Composition

	Air	Water	Sucrose Solution		
Sucrose Fraction $(w_s/w_w)$	0 %	0 %	50~%	66.66~%	100~%
Sugar Mass (kg)	0	0	1495.5	1993.8	2991
Fluid Mass (kg)	0	0	2991	2991	2991
Density $(kg m^{-3})$	1.225	997	1138.5	1175.62	1224.25
Viscosity $(mPa \cdot s)$	0.016	1	3.49	5.17	11.5



## Impact of simulation order



## Observed regimes

### **Rolling Regime**





### **Cascading Regime**





## Observed regimes

### **Cataracting Regime**



### **Centrifuging Regime**







# Regime capture efficiency

## Methodology:

- Calibrate friction coefficients in rolling regime for a given fluid
- Simulate all velocity range and compare the regime obtained in the simulation with reality

## Result for accuracy:

 96% 3mm glass, 84% for 6mm ABS, <u>54% for 5mm</u> <u>glass</u>



# Interesting particle rotation







# Challenge and future direction

Onwards to Matrix-free Methods



# What are the challenges?

Time consuming part matrix-based algorithm

- Assembly of the system matrix
- Assembling the preconditioner (e.g. AMG)
- Solving the linear system

## Motivation for a matrix-free approach

- On modern CPUs access to main memory is the bottleneck
- Krylov methods (e.g. GMRES) use matrix-vector multiplications
- For large problems, it is faster to compute matrix entries than to load them



## Matrix-free: Main idea

Calculate the matrix-vector multiplication defining an operator



Other elements

- Geometric multigrid preconditioning with global coarsening
- Other optimization (vectorization and sum-factorization)



# Result for Taylor-Green vortex



As soon as go to Q2, matrix-free outperforms matrix-based





## Conclusions





# Challenges and conclusion

## Modeling granular flow remains highly challenging

- Lagrangian strategies are computationally intensive
- Turbulent particle-fluid interaction and turbulence modeling in the presence of particles is difficult

## High-order CFD-DEM are part of the solution

- Combined order and mesh resolution give lots of freedom
- Matrix-free is the way to go to reduce computational cost
- Open-source computational model are a part of the solution • Reproducibility
- Capacity to be launched at large scale
- Imperfect, but that's all right







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## Thanks! Bruno Blais

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## Wish to learn more?

- •El Geitani, T., Golshan, S., & Blais, B. (2023). Toward High-Order CFD-DEM: Development and Validation. Industrial & Engineering Chemistry Research, 62(2), 1141-1159.
- •El Geitani, T., & Blais, B. (2023). Quadrature-Centered Averaging Scheme for Accurate and Continuous Void Fraction Calculation in Computational Fluid Dynamics–Discrete Element Method Simulations. Industrial & Engineering Chemistry Research, 62(12), 5394-5407.
- •Ferreira, V. O., El Geitani, T., Junior, D. S., Blais, B., & Lopes, G. C. (2023). In-depth validation of unresolved CFD-DEM simulations of liquid fluidized beds. *Powder Technology*, 118652.
- •Geitani, TE, Golshan, S, Blais, B. A high-order stabilized solver for the volume averaged Navier-Stokes equations. Int J Numer Meth Fluids. 2023; 95( 6): 1011–1033.
- •Golshan, S., Munch, P., Gassmöller, R., Kronbichler, M., & Blais, B. (2023). Lethe-DEM: An open-source parallel discrete element solver with load balancing. *Computational Particle Mechanics*, *10*(1), 77-96.
- Barbeau, Lucka, et al. "High-order moving immersed boundary and its application to a resolved CFD-DEM model." Computers & Fluids 268 (2024): 106094.

Barbeau, Lucka et al. In preparation (2024)

