

# Direct numerical simulation of a hypersonic boundary layer with porous wall injection

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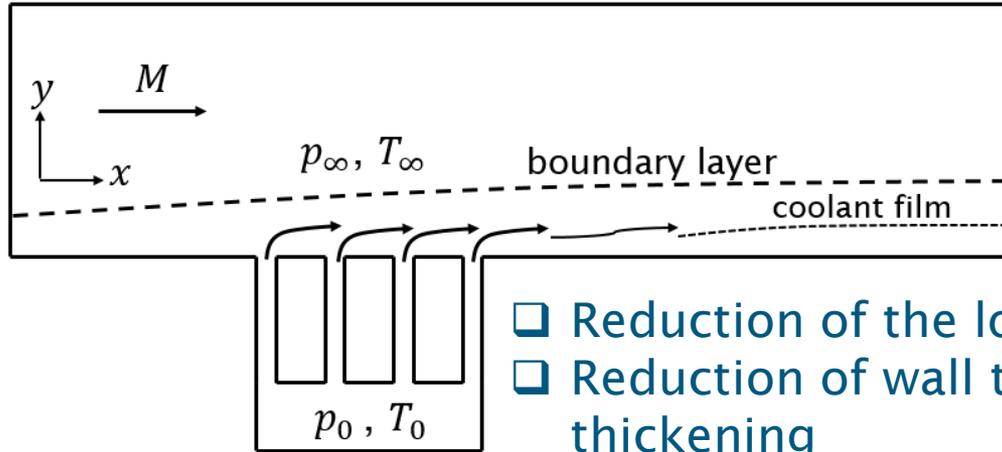
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Part of EPSRC Programme Grant *Transpiration Cooling Systems for Jet Engine Turbines and Hypersonic Flight*  
<http://transpirationcooling.eng.ox.ac.uk>



# Film cooling



- ❑ Formation of a **thin film of cold fluid adjacent** to the wall
- ❑ Wall heat flux reduction achieved by a dual effect:

- ❑ Reduction of the local fluid temperature near the wall
- ❑ Reduction of wall temperature gradient due to BL thickening

Film cooling (Fitt et al., 1985, 1994): effusion cooling vs transpiration cooling

- ❑ **Effusion cooling** (Wittig et al. 1996, Baldauf et al. 2001): coolant injected by localized holes

- 3D effects induced in the BL, which promote transition
- Typically used on gas-turbine blades (turbulent BL)

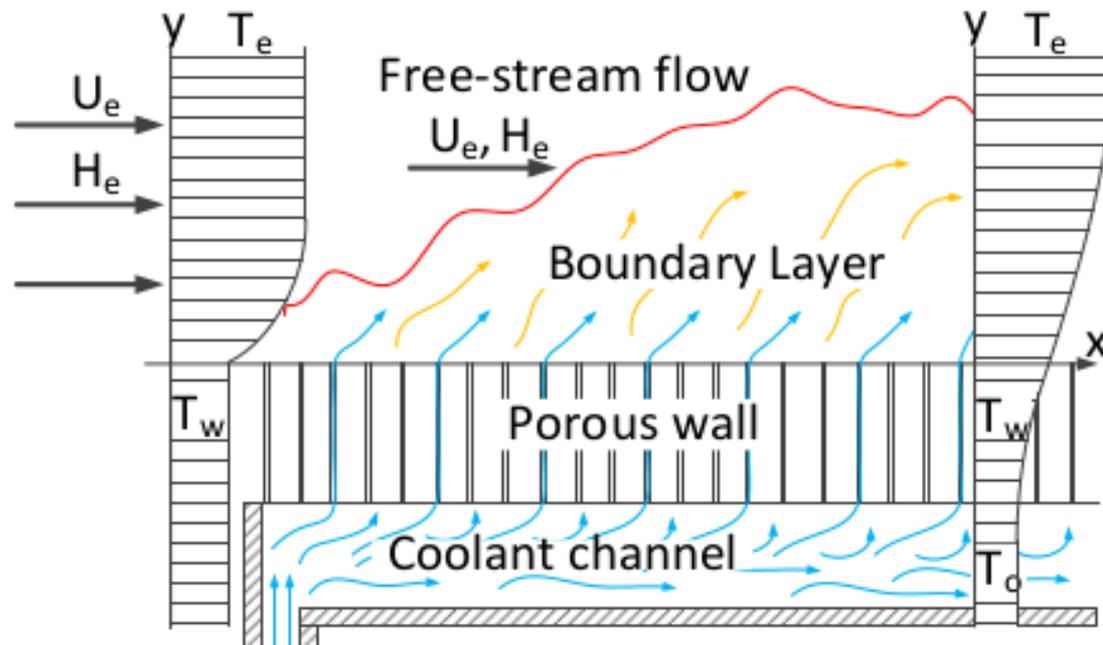
- ❑ **Transpiration cooling** (Meinert et al. 2001, Langener et al. 2011): fluid injected through a porous material

- More homogeneous injection
- Higher effectiveness



# Purpose of direct numerical simulation

- Understand the physics of blowing in hypersonic flow through porous surfaces
- Predict cooling performance – heat transfer rates
- Estimate the effects on BL stability and transition
- Validate results with experiments



# Numerical method for full Navier-Stokes equations

- **Finite volume** approach in a Cartesian reference system with embedded boundaries
- Structured-adaptive-mesh-refinement (**SAMR**) algorithm (AMROC - Deiterding, 2005 & 2011): the Cartesian grid is **locally refined** by adding consecutive finer grid levels during the iteration cycles (patch-wise refinement strategy)
- **Hybrid WENO-CD** method up to **6<sup>th</sup> -order** accurate for both inviscid and viscous fluxes (Cerminara, Deiterding, Sandham, 2018)
- Central Differencing (CD) scheme used in the smooth regions:  
E.g. for 6<sup>th</sup> -order, the numerical flux at the interfaces is evaluated as

$$\hat{f}_{i+\frac{1}{2}} = \alpha(f_{i+3} + f_{i-2}) + \beta(f_{i+2} + f_{i-1}) + \gamma(f_i + f_{i+1})$$

$$\alpha = 1/60 \quad \beta = -2/15 \quad \gamma = 37/60$$

- 3<sup>rd</sup> -order Runge-Kutta for time integration
- **Characteristic-based switch** to a high-resolution WENO scheme to handle discontinuities in a computationally efficient way (Hill and Pullin, 2004)



# Shock-capturing scheme

- Weighted-essentially-non-oscillatory - symmetric - order-optimized (**WENO-SYMOO**) scheme (Martin et al. 2006)
- **Reduced dissipation** (compared to 5th-order WENO-JS, Jiang and Shu 1996), and **6th** maximum formal **order** of accuracy, reached through:
  1. Symmetrization: a **fourth** candidate **stencil** is **added** to the initial ( $r=3$ )-points upwinded-biased candidate stencils of the WENO-JS
  2. Optimal weights  $C_k$  to guarantee 6th-order accuracy

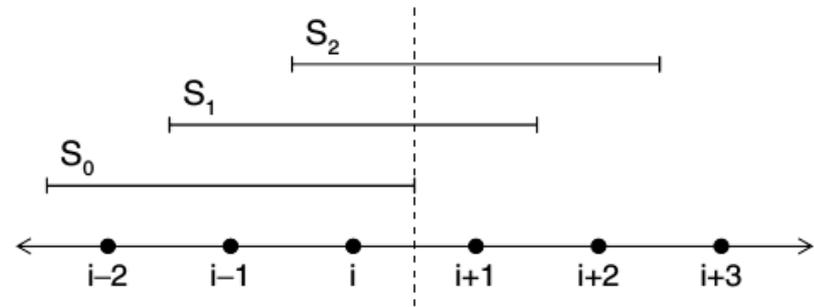
$$\hat{f}_{i+\frac{1}{2}}^+ = \sum_{k=0}^r \omega_k q_k^r \quad q_k^r = \sum_{l=0}^{r-1} d_{k,l}^r f(u_{i-r+k+l+1})$$

r-th-order polynomial interpolation on the k-th stencil

$$\omega_k = \frac{\alpha_k}{\sum_{k=0}^{r-1} \alpha_k} \quad \text{weight of the k-th interp.} \quad \alpha_k = \frac{C_k}{(\epsilon + IS_k)^p}$$

$$IS_k = \sum_{m=1}^{r-1} \left( \sum_{l=0}^{r-1} b_{k,m,l}^r f(u_{i-r+k+l+1}) \right)^2$$

smoothness indicator



from Martin et al. 2006



# Scheme switch

- Riemann problem is solved at each cell interface through Roe's averaged state
- Lax condition is applied to detect shock/rarefaction waves

$$|u_R \pm a_R| < |u^* \pm a^*| < |u_L \pm a_L|$$

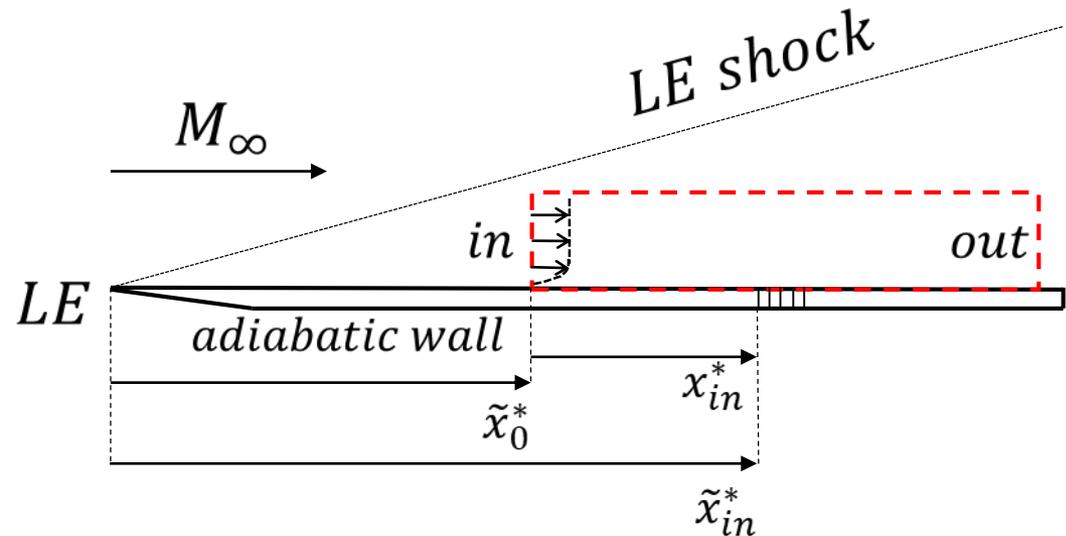
- A threshold is applied to the residuals between the left/right state and the intermediate state to select the strong waves and neglect the weak ones (Ziegler et al., 2011)
- Another threshold is applied to the function  $\phi(\theta_i)$  to allow WENO only in the high pressure-gradient regions (Ziegler et al., 2011):

$$\phi(\theta_i) = \frac{2\theta_i}{(1 + \theta_i)^2} \quad \theta_i = \frac{|p_{i+1} - p_i|}{|p_{i+1} + p_i|}$$



# Flat plate studies: Flow conditions and domain set-up

- $M_\infty = 5$
- $Re_m = 12.6 \times 10^6 \text{ 1/m}$
- $T_\infty^* = 81.7 \text{ K}$
- $T_{0\infty}^* = 490.2 \text{ K}$
- $U_\infty^* = 906 \text{ m/s}$
- $\rho_\infty^* = 0.078 \text{ Kg/m}^3$
- $p_\infty^* = 1.832 \times 10^3 \text{ Pa}$
- $\rho_\infty^* U_\infty^{*2} = 6.4 \times 10^4 \text{ Pa}$
- $\mu_\infty^* = 5.6 \times 10^{-6} \text{ Kg/(m} \cdot \text{s)}$
- $T_w^* = T_{ad}^* = 5.24 T_\infty^*$
- $\tilde{x}_0^* = 127 \text{ mm}$
- $\tilde{x}_{in}^* = 182 \text{ mm}$
- $x_{in}^* = 55 \text{ mm}$



$\tilde{x}^*$  = distance from LE

$x^*$  = distance from inflow in the domain coordinate system



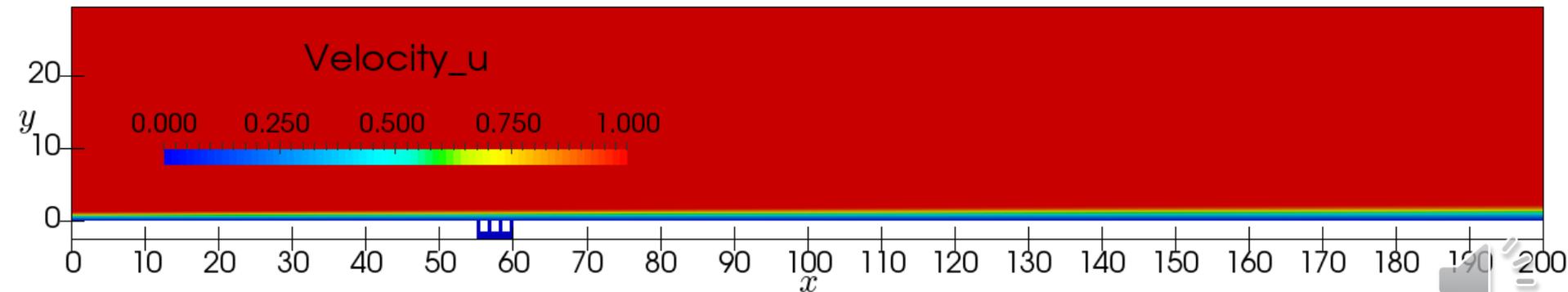
# Initialization

From the inflow profile, the displacement thickness of the laminar BL in the initial solution grows in the streamwise direction, following the relation:

$$\frac{\delta^*(\tilde{x}^*)}{\delta_0^*} = \Delta \frac{\sqrt{2Re_{\tilde{x}^*}}}{Re_{\delta_0^*}},$$

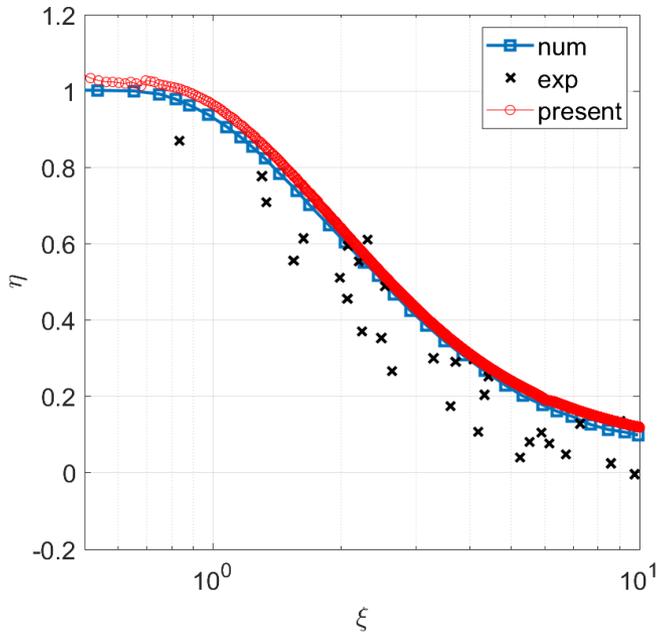
where  $\Delta = 8.18$  (for the present case) is a scaling factor from the similarity solution.

The result for the initial state at all the  $x$  positions is:

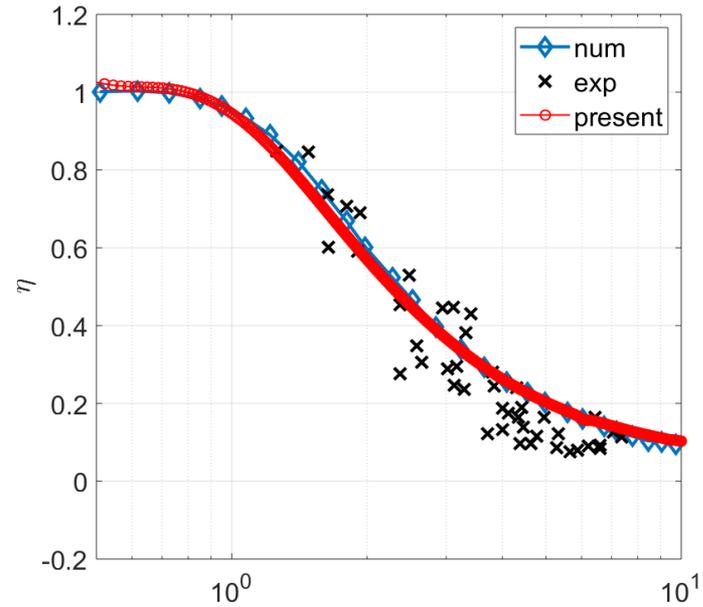


# Validation case: slot injection with multi-component gas effects

## Air injection

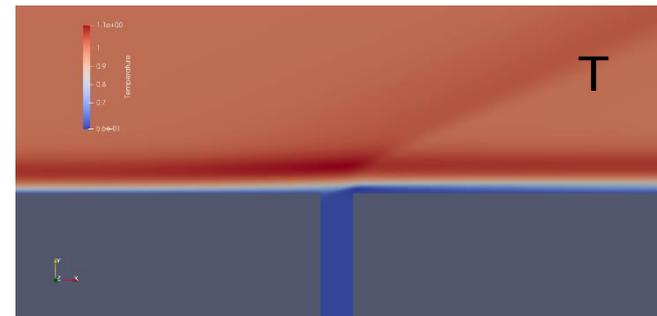


## CO<sub>2</sub> injection



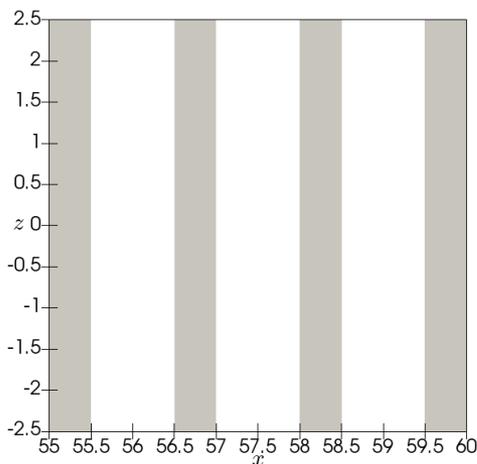
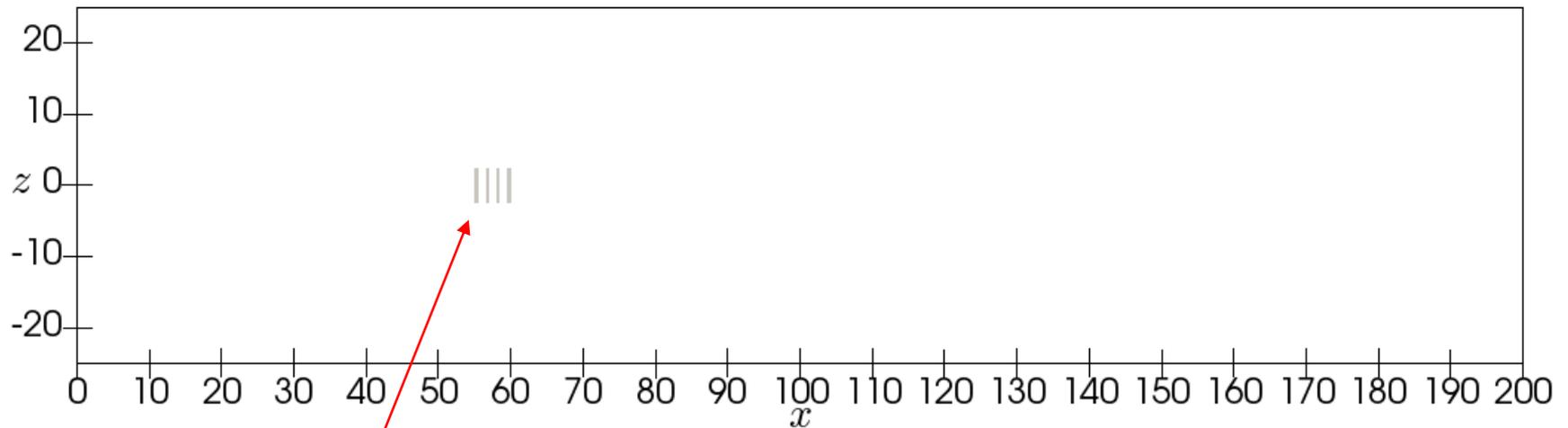
$$\eta = 1 - \frac{q_{w,c}}{q_{w,nc}}$$

Very good agreement with Keller et al. (2015) results of cooling effectiveness in a Mach 2.6 case



# 3D flat plate with blowing through finite slots

Grid size (base level)  $N_x \times N_y \times N_z = 800 \times 384 \times 100$



Domain size  $L_x \times L_y \times L_z = 200 \times 32 \times 50$

Width of the slot region = 5 mm

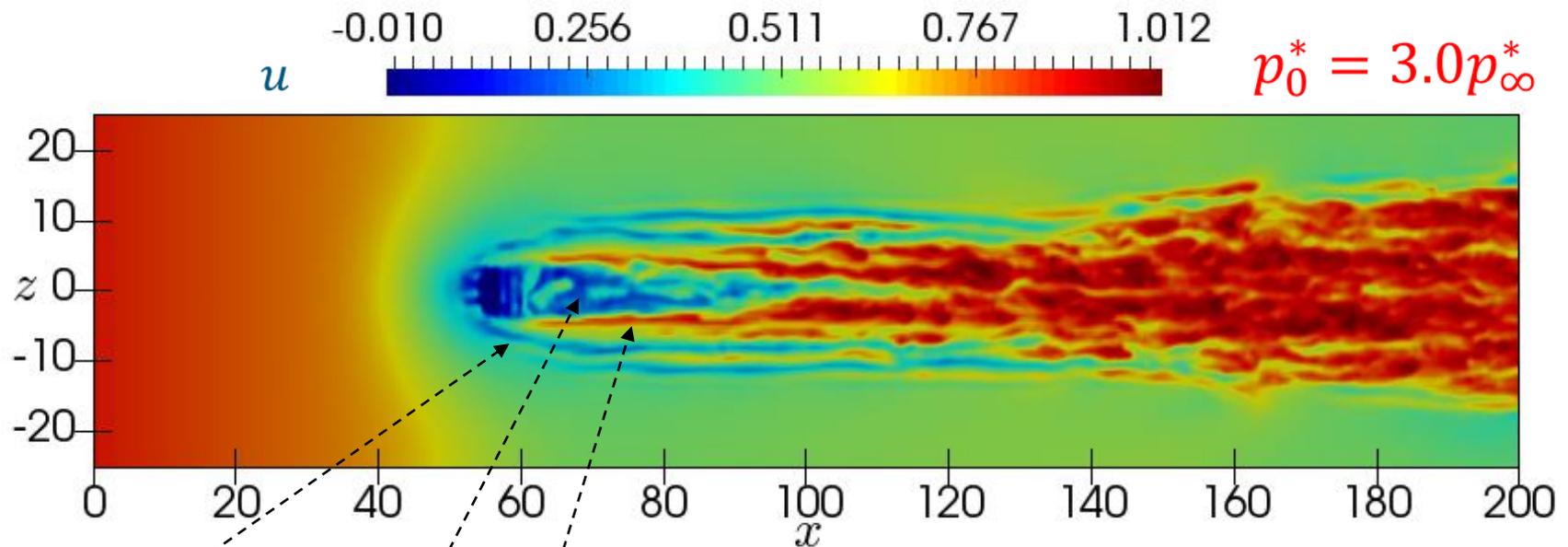
Minimum cell sizes  $\Delta x = 0.125$  mm,  $\Delta y = 0.0415$  mm,  $\Delta z = 0.25$  mm (with 2 grid levels)

Simulation run with 480 cores on Iridis 4 cluster (Soton HPC facility)



# 3D flat plate with blowing through slots

$$y = 1 \text{ (} y^* = 1 \text{ mm)}$$



Horseshoe vortices (low-velocity streaks)

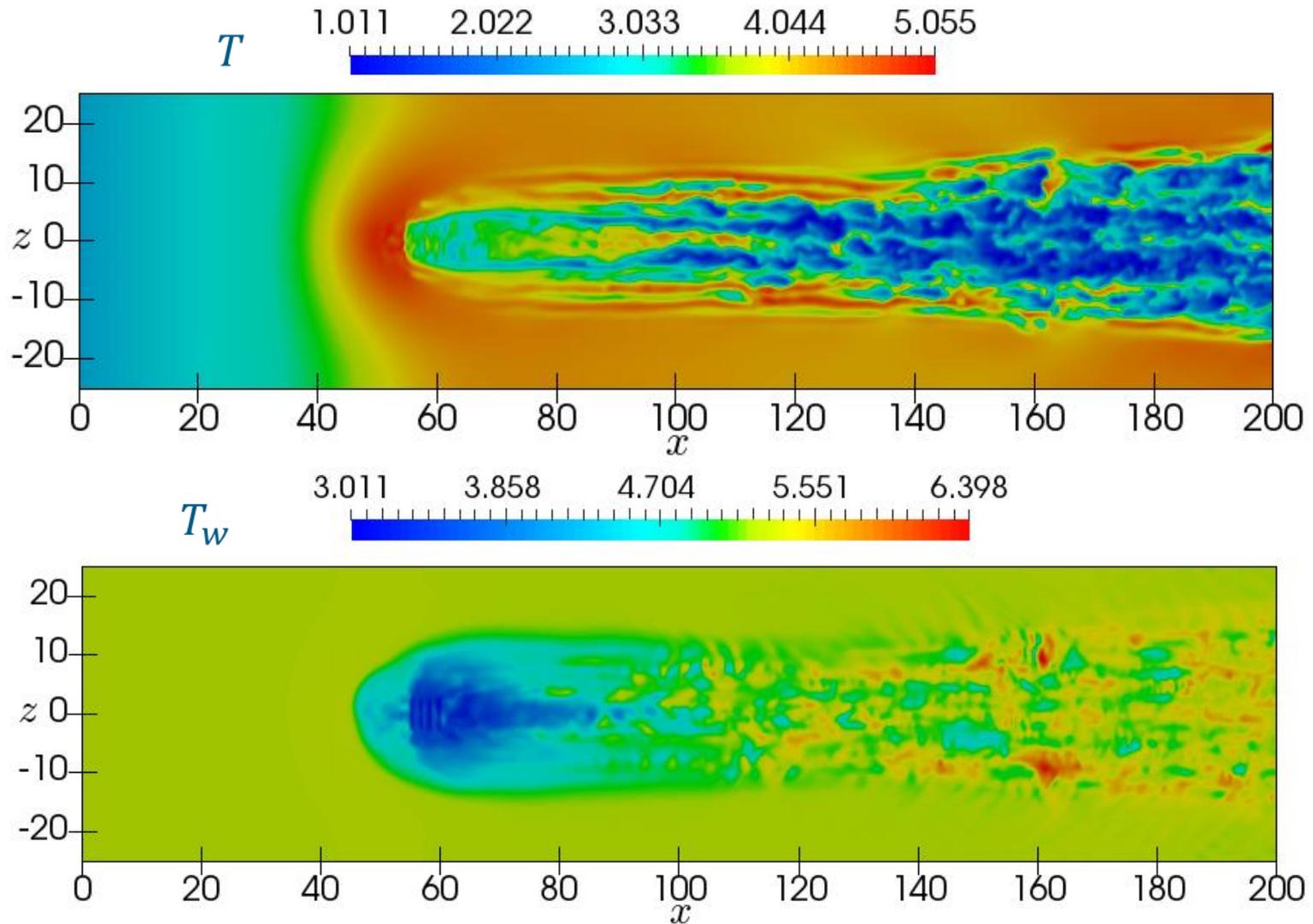
wake

Edge vortices (high-velocity streaks)

- Transition induced by breakdown of the main edge vortices downstream of  $x = 90$  ( $x^* = 90 \text{ mm}$ ,  $\tilde{x}^* = 217 \text{ mm}$ )
- Wedge-shaped transition front spreading downstream



# 3D flat plate with blowing through slots



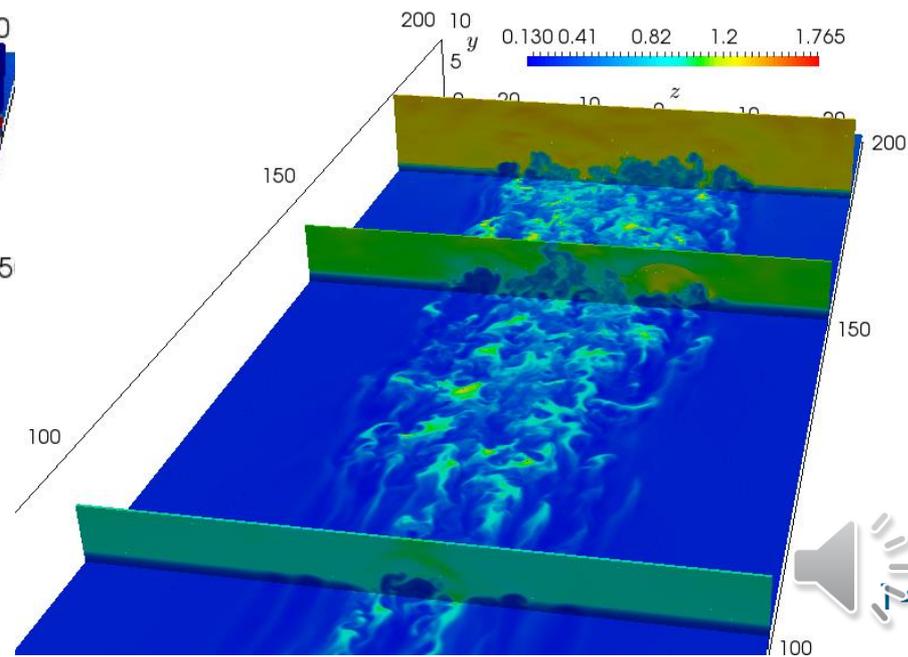
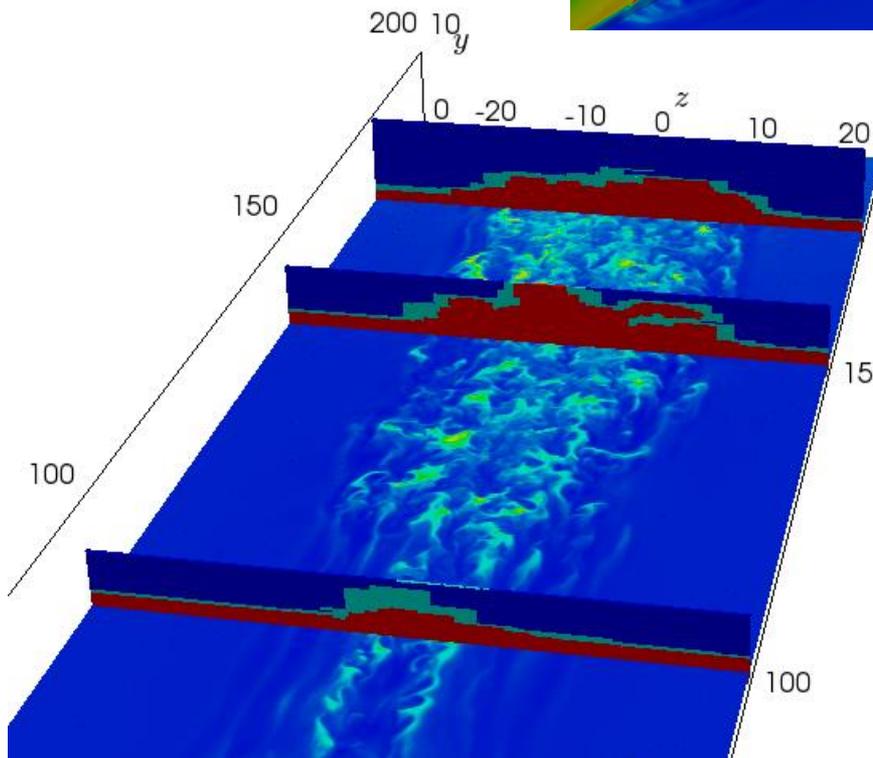
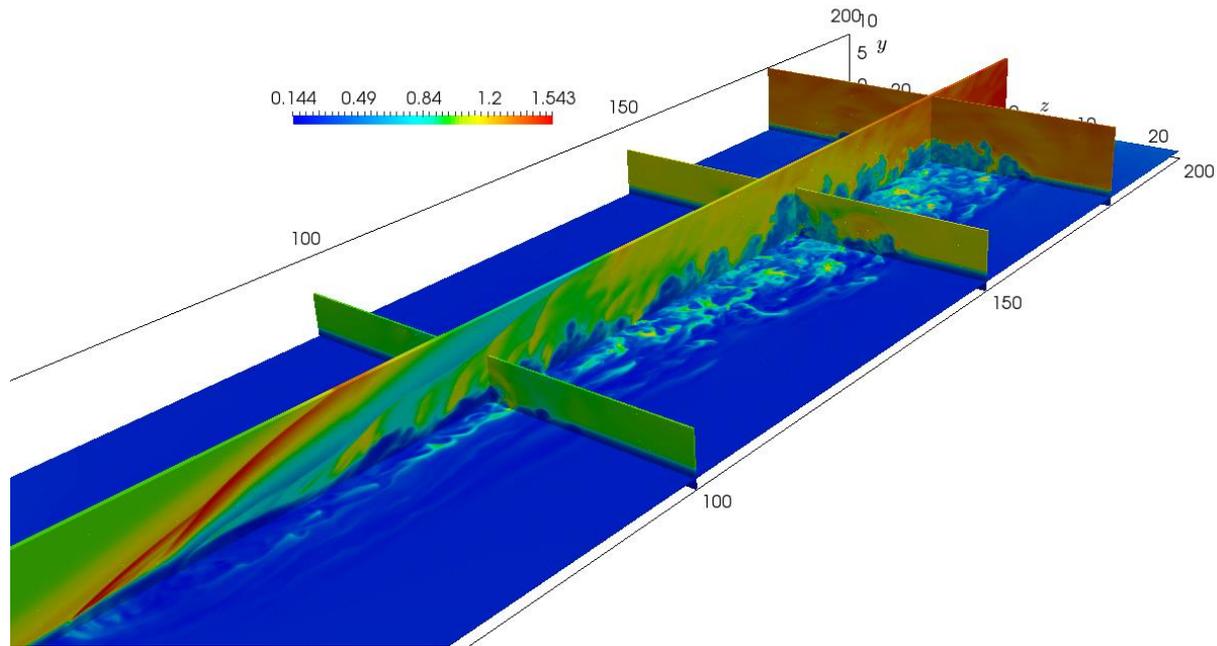
- Cooling performed mainly at the slot sides and along the wake
- Limited streamwise extent of the cooled region due to transition to turbulence



# 3D results

Right, bottom right: density field

Left, bottom left: AMR levels

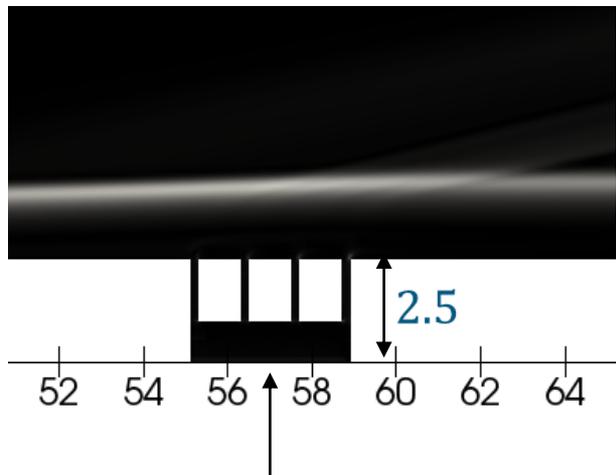
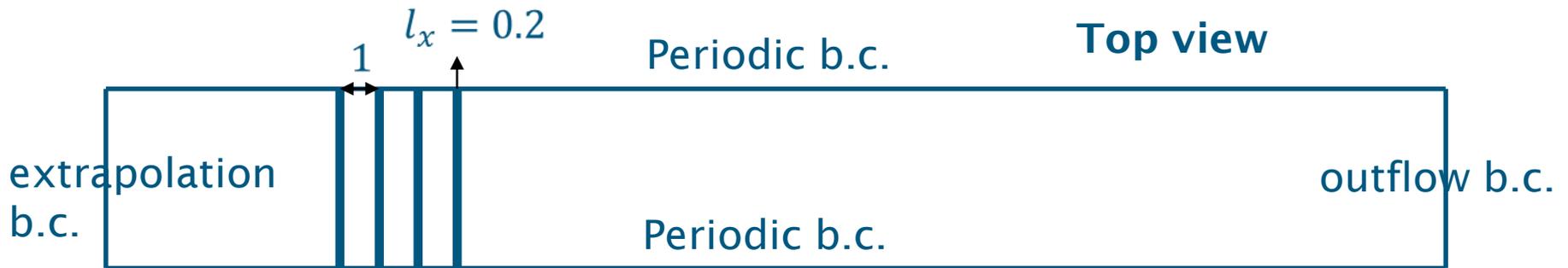


# Oxford slot experiments

Characteristic length = 1 mm (BL displacement thickness at inflow)

$$L_x = 160, L_y = 32, L_z = 4$$

$$N_x \times N_y \times N_z = 3200 \times 384 \times 40$$



Plenum b.c. :  $p_0, T_0 = 290 K$

Injection conditions:

1. *no injection*

2.  $\frac{p_0}{p_\infty} = 1.6$

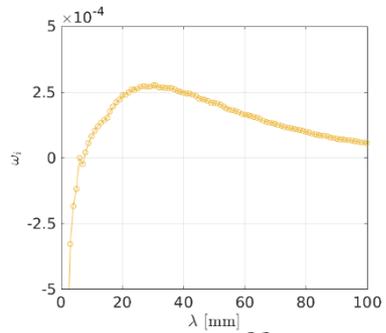
3.  $\frac{p_0}{p_\infty} = 1.8$

4.  $\frac{p_0}{p_\infty} = 2.7$

5.  $\frac{p_0}{p_\infty} = 4.8$



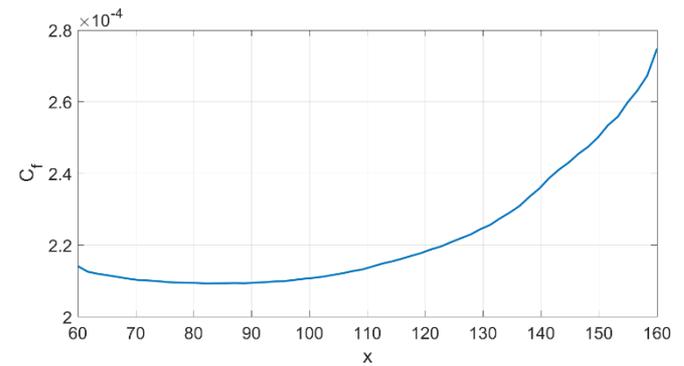
# Transitional flow state



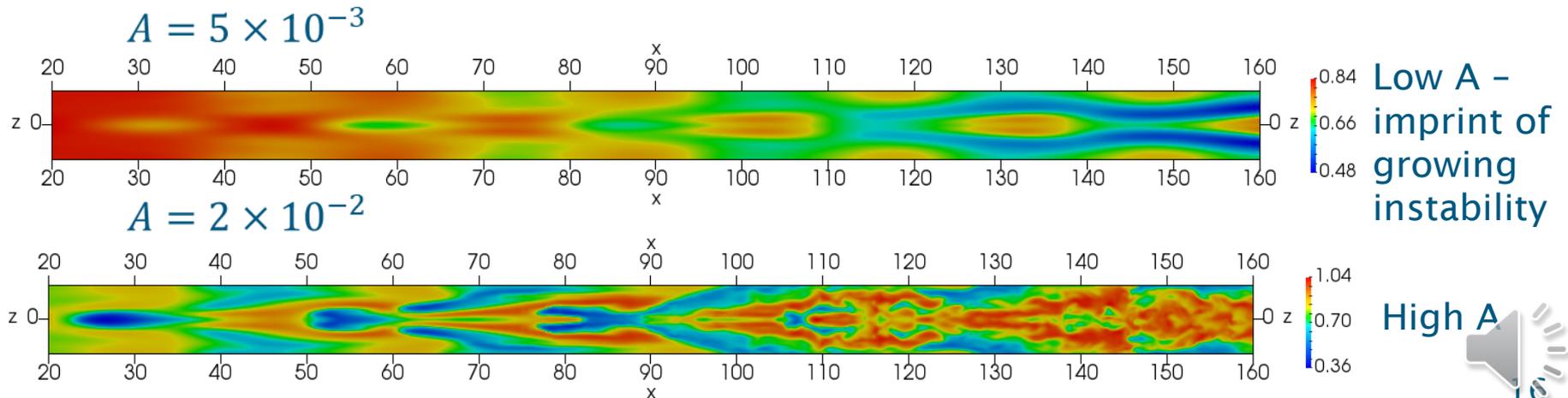
LST to identify the most unstable frequencies and wavenumbers

Model to reproduce instability modes

$$v'(x, z, t) = \sum_{m=1}^M A \cos(\beta_m z + \phi_m) \cos[\alpha(x - x_0) - \omega t + \psi]$$

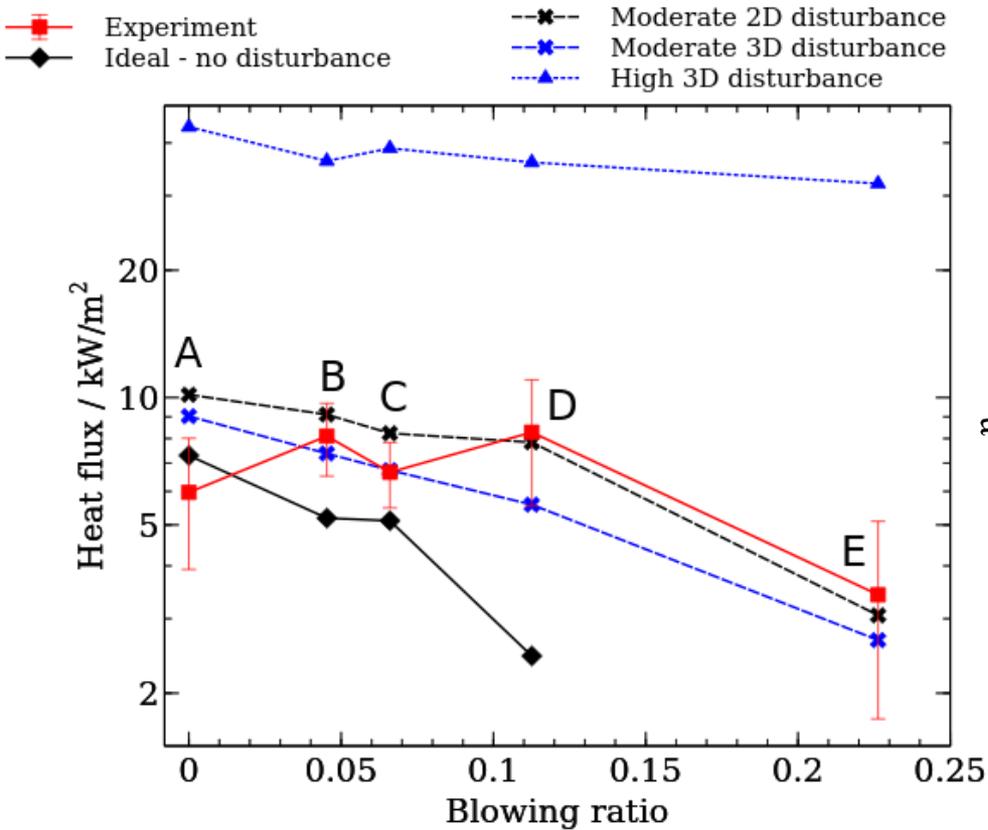


Most unstable mode:  $\alpha=0.2$ ,  $\beta=0.78$

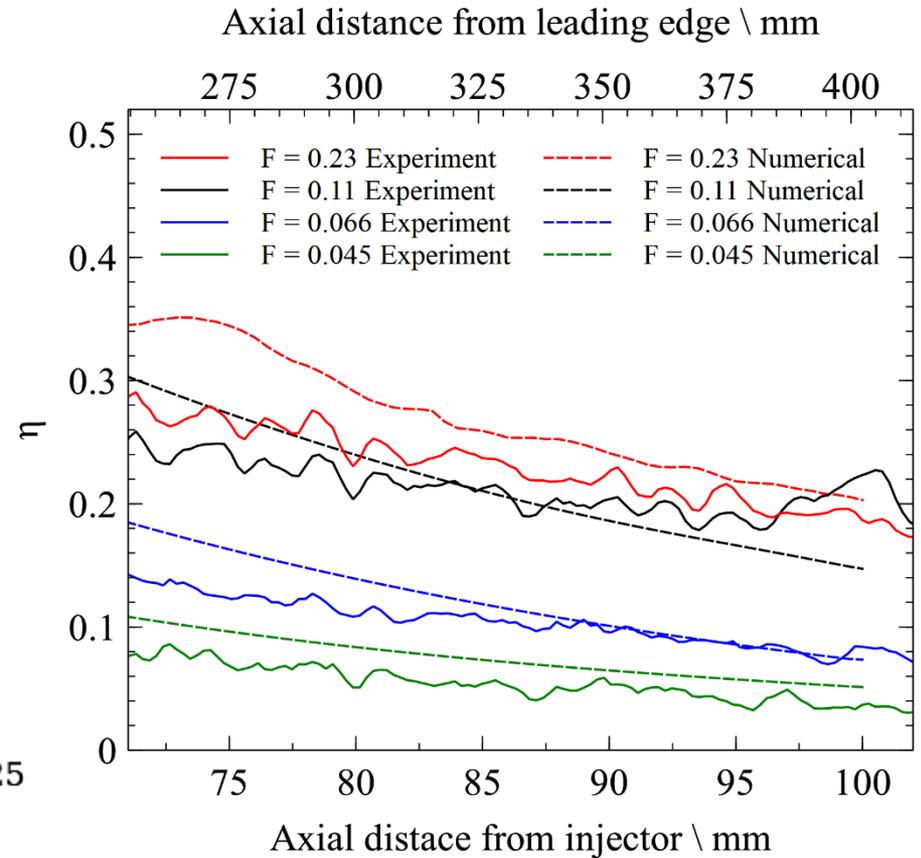


# Heat flux and effectiveness predictions

## Heat flux



## Effectiveness measure based on concentration



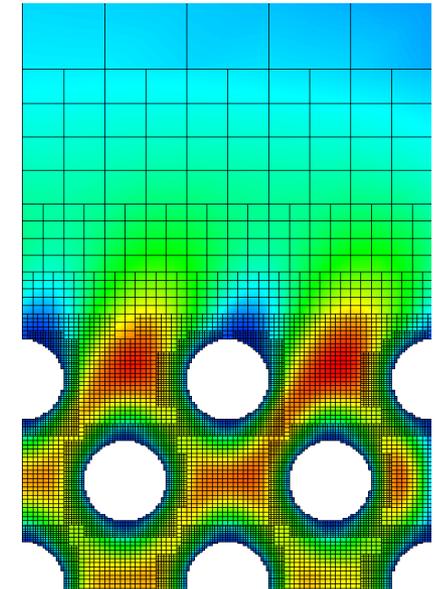
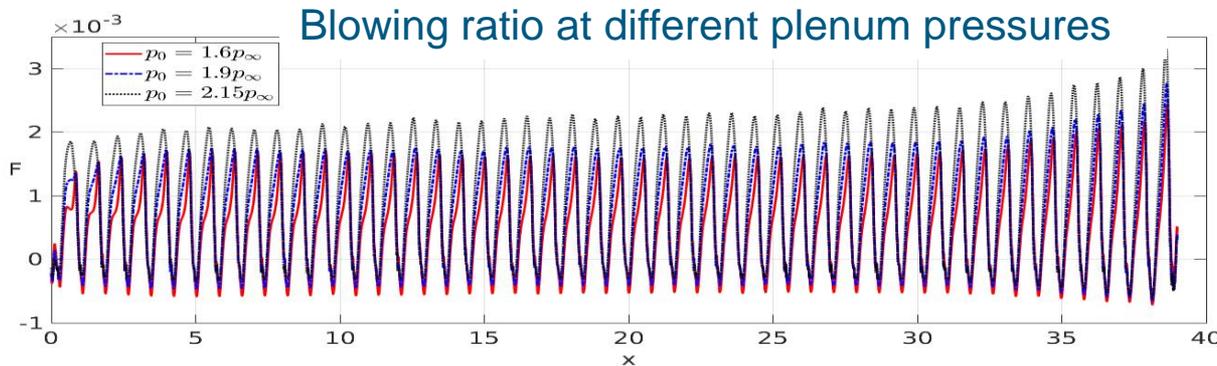
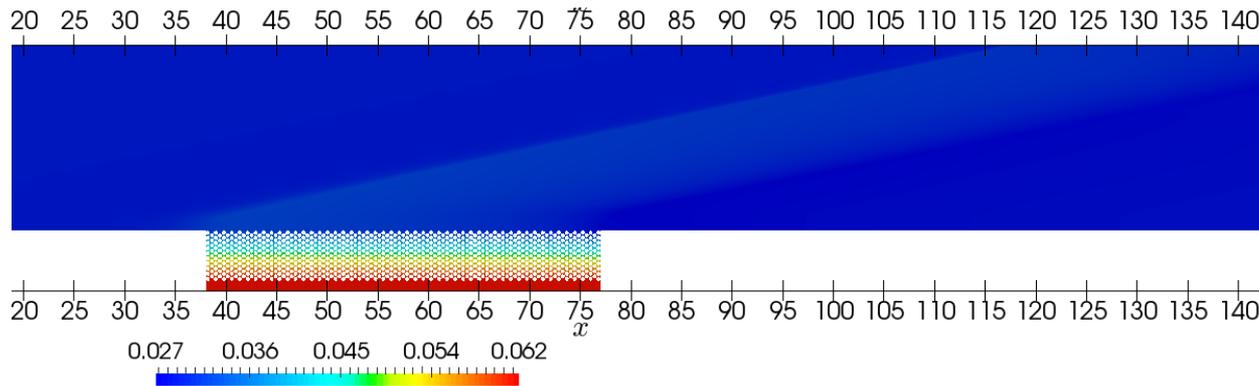
(Cerminara et al., Aerospace Science Technology, 2020)



# Modelling of porous material

- Original project proposal assumed grain structures of  $>50\mu\text{m}$
- Used material has grains  $<2\mu\text{m}$  (cost of DNS in porous material increases by factor approx.  $25^4=390625$ )
- Our approach: model the porous media flow using regular sphere/cylinder arrays of resolvable mesoscopic size

Cylinders with radius of  $200\mu\text{m}$ , pressure contours



Automatic adaptive mesh (vertical velocity)

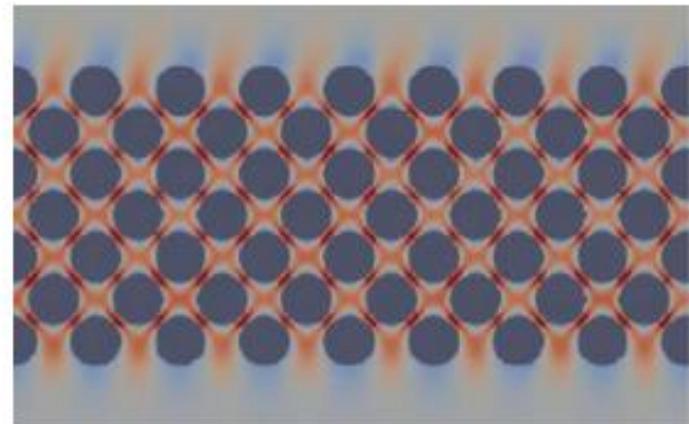
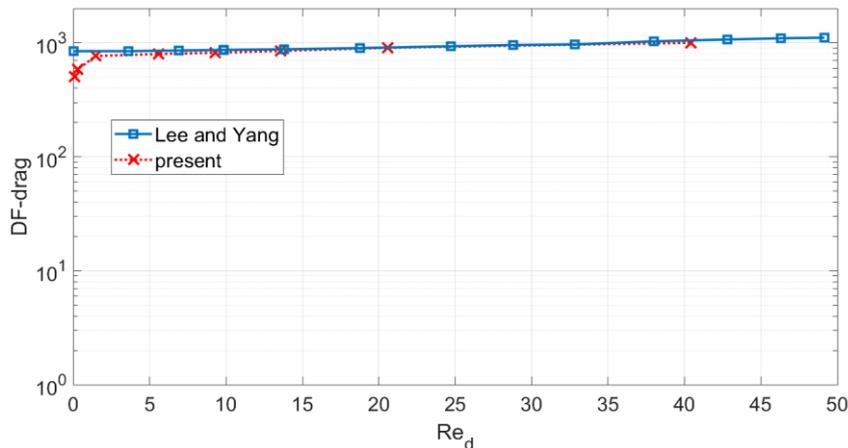
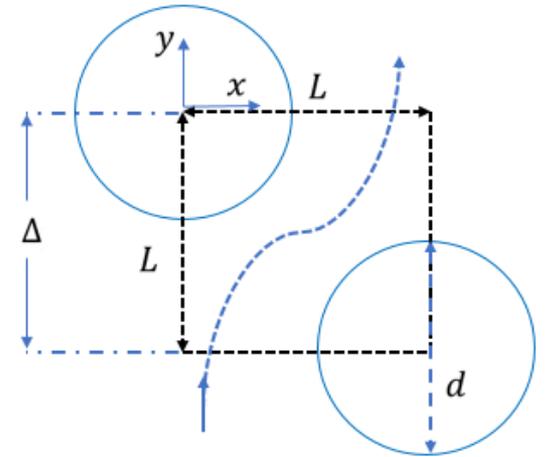


# Mesososcopic Darcy-Forchheimer model

$$\left(\frac{d^2}{K} + FRe_d\right) = \frac{\Delta p}{L} \frac{d^2}{\mu U} = D - F \text{ drag} \quad Re_c = \frac{\rho U_c L}{\mu} \quad U_c = \sqrt{\frac{\Delta p}{\rho}} \quad Re_d = \frac{qd}{\mu} = Re_c \left(\frac{d}{L}\right) q$$

$$q = \frac{1}{\rho U_c} \frac{1}{L} \int_0^L \rho u dx \quad \frac{d}{L} = f(\varepsilon) \quad \Delta p = \rho \left(\frac{Re_c \mu}{\rho L}\right)^2$$

- Objective: represent the pressure drop of micro-structured material with larger cylinders
- Idea is to vary  $d$  but maintain  $q$
- Preliminary study follows approach of Lee and Yang (1997)



Zoom in: y-momentum in porous layer for  $d=12\mu\text{m}$



# Matching with experiment in terms of blowing ratio, and rescaling to higher pore sizes

Quantifying integrated blowing ratio solution ( $q$ ) for the different pore/particle sizes:

$$Re_C = 1.67$$

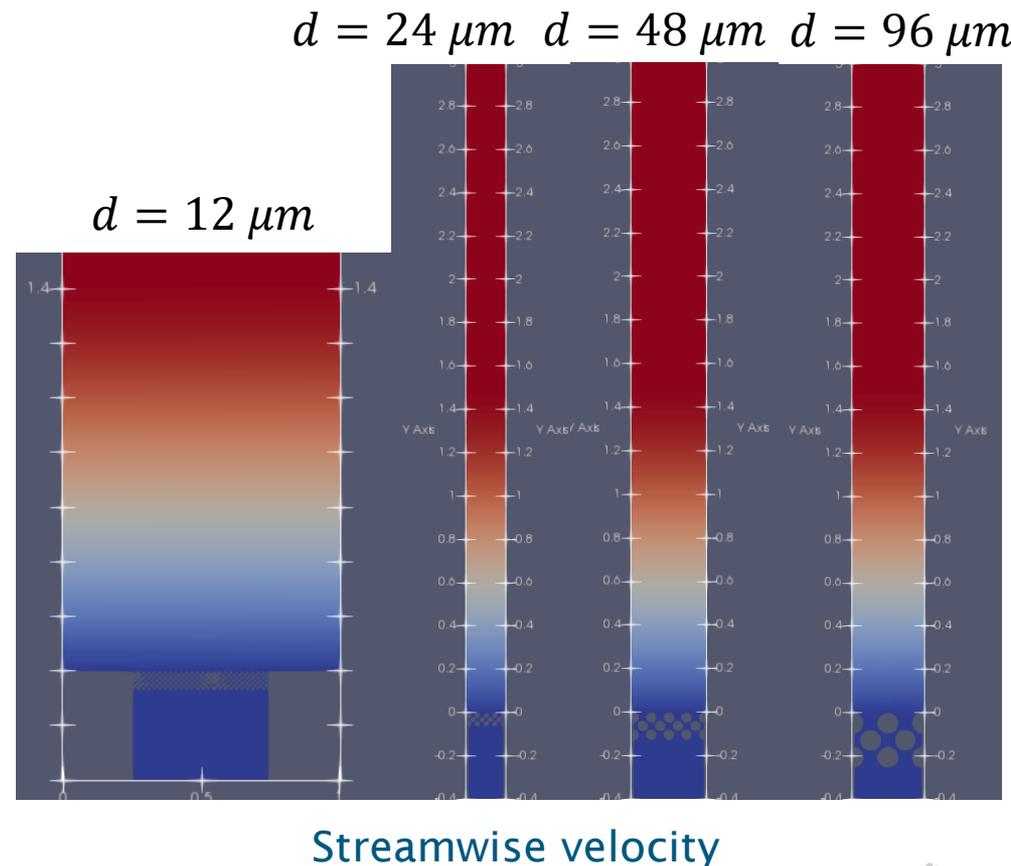
$$d = 12 \mu\text{m} \rightarrow q = 3.8 \times 10^{-3}$$

$$d = 24 \mu\text{m} \rightarrow q = 3.7 \times 10^{-3}$$

$$d = 48 \mu\text{m} \rightarrow q = 3.6 \times 10^{-3}$$

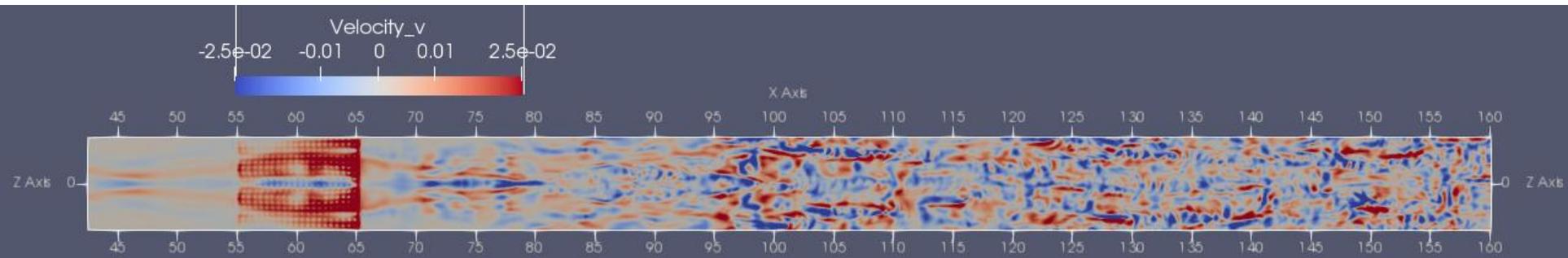
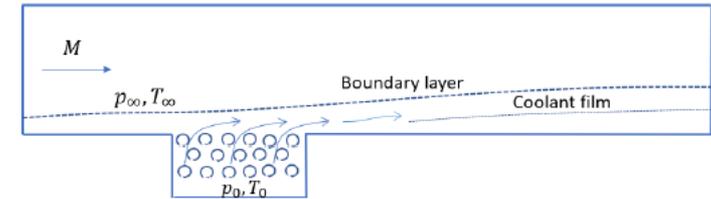
$$d = 96 \mu\text{m} \rightarrow q = 3.8 \times 10^{-3}$$

Simulation results show rescaling capability from very small pore sizes to higher scales in terms of the D-F behaviour

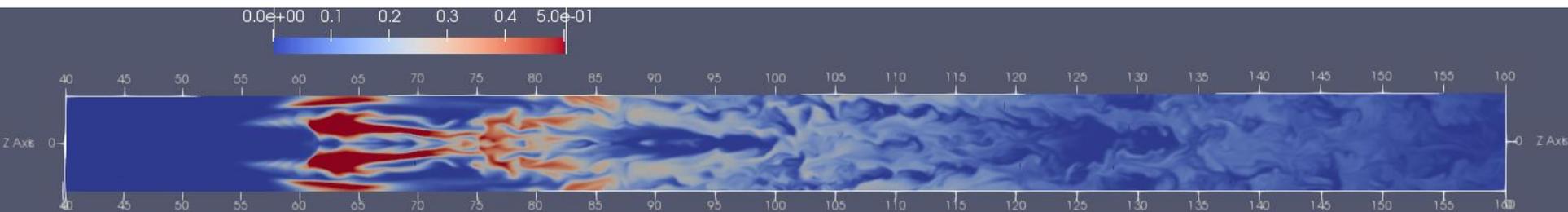


# Oxford ultra-high temperature ceramics experiment

- $M=5$ ,  $Re=12.6E6$  1/m from Herman et al. HiSST (2018)
- $F_{porous}=0.002$ , porosity 42% modelled with spheres of diameter  $d=340\mu m$
- $p_0=1.3 p_{inf}$  to match blowing ratio
- 2000 x 300 x 100 mesh plus 2 levels refined factor 2



v-velocity: note specimen on left and transition to turbulence towards the right

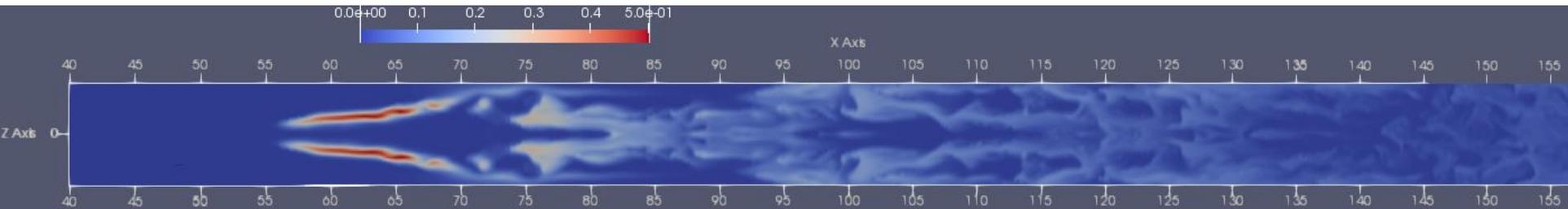


Coolant concentration from top and side view

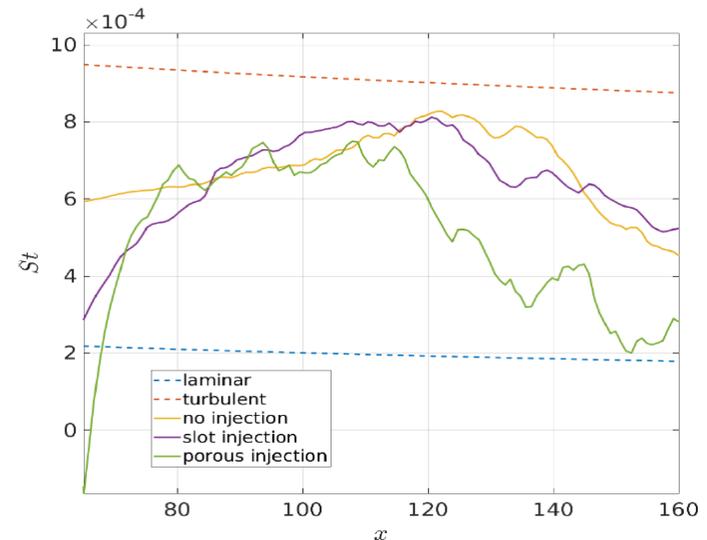
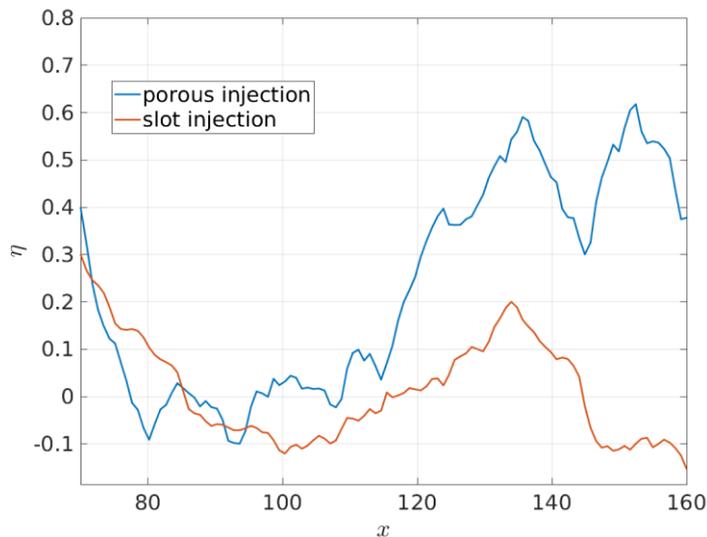


# Comparison with suitable slot case

- Chose a slot case with comparable non-dimensional mass flow of injected coolant  $d = \rho U A_{inj}$
- $F_{porous} = 0.002$ ,  $F_{slots} = 0.045$ ,
- $d_{slots} = 0.173$ ,  $d_{porous} = 0.144$



Cooling effectiveness (averaged spanwise and in time)      Stanton number (averaged spanwise and in time)



(Cerminara, Deiterding, Sandham. Int. J. Heat Fluid Flow, 2020)



# Conclusions

- A hybrid 6<sup>th</sup> order multi-species WENO/CD method in the AMROC framework for DNS of hypersonic boundary layers with wall injection has been implemented
  - Method is efficient, robust and runs well on large parallel systems
- Have developed a novel mesoscopic model for flow through porous structures that can be used with accurate hypersonic boundary layer DNS, thereby avoiding resolving the  $O(1-2\mu\text{m})$  structures of the manufactured material
- Validation cases for both slot injection and porous injection have demonstrated a good agreement with available experimental results
- Level of realism in order to model realistic flight conditions will be increased further by supplementing the implementation with
  - Transient surface heat conduction modelling
  - Thermochemical and vibrational nonequilibrium models